

# Time evolution:

Coupled ODE's: (discretized PDE's  
for instance !)

$$\dot{z} = f(z)$$

$$\vec{z} = \begin{pmatrix} z_1 \\ \vdots \\ z_n \end{pmatrix}$$

Aim: Given  $z(t=0)$ , find  $z(t)$ .

Method of choice: timestepping

$$\vec{z}(t_1) = M(\vec{z}(t_0))$$

$$\vec{z}(t_2) = M(\vec{z}(t_1))$$

...

$M$ : Some map (algorithm!)

Today: learn from examples

① Orbit in a radial potential



$$V(r) = -\frac{\lambda}{\sqrt{x^2 + y^2}}$$

$$E.o.M: \dot{p}_x = -\frac{\partial V}{\partial x} = -\frac{\lambda x}{\sqrt{x^2 + y^2}} \equiv K_x(x, y)$$

$$\dot{p}_y = -\frac{\partial V}{\partial y} = -\frac{\lambda y}{\sqrt{x^2 + y^2}} \equiv K_y(x, y)$$

$$\dot{x} = p_x$$

$$\dot{y} = p_y$$

$$H = \frac{p_x^2}{2} + \frac{p_y^2}{2} - \frac{\lambda}{\sqrt{x^2 + y^2}}$$

$$L = x p_y - y p_x$$

Energy and angular momentum are conserved  $\Rightarrow$

$$r(\theta) = \frac{-L^2}{1 + e \cos(\theta)}, \quad e = \sqrt{1 + 2EL^2}$$

Attempt  $\times 1$ : "Forward Euler" (FE)

$$\frac{df}{dt} \approx \frac{f(t+\Delta t) - f(t)}{\Delta t}$$

$$\Rightarrow \vec{p}(x+\Delta t) = \vec{p}(t) + \Delta t \vec{K}(x(t), y(t))$$

$$\vec{x}(x+\Delta t) = \vec{x}(t) + \Delta t \vec{p}(t)$$

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FE does not preserve L cons.

Notation:  $f(t+\Delta t) \equiv f^{+1}$   $f(t) \equiv f$

$$K(x(t+\Delta t), y(t+\Delta t)) = K^{+1}$$

$$L^{+1} = x^{+1} p_y^{+1} - y^{+1} p_x^{+1}$$

$$x^{+1} p_y^{+1} = x p_y + \Delta t p_x p_y + \Delta t K_y x + \Delta t^2 K_y p_x$$

$$\Rightarrow L^{+1} = L + \Delta t (\underbrace{K_y x - K_x y}_{=0}) + \Delta t^2 (K_y p_x - K_x p_y) + \Delta t^2 (K_x p_y - K_y p_x)$$

Suggests:  $\begin{cases} \vec{p}^{+1} = \vec{p} + \Delta t \vec{K} \\ \vec{x}^{+1} = \vec{x} + \Delta t \vec{p}^{+1} \end{cases}$

Semi-implicit Euler

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Take home: Physics  
↑ ↓  
algorithms

Built an algorithm that  
preserves the conservation laws  
("symplectic")

↓  
Good for Hamiltonian systems

$$H: H(x, p_x) \{x, p\} = 1$$

$$\text{Poisson bracket: } \{f, g\} = \frac{\partial f}{\partial x} \frac{\partial g}{\partial p_x} - \frac{\partial f}{\partial p_x} \frac{\partial g}{\partial x}$$

Symplectic: Preserves  $\{\cdot, \cdot\}$



Preserves conservation laws

Some algorithm's properties:

- Explicit vs implicit
  - + cheaper
  - less stable
  - More expensive
  - + More stable
- Order of convergence
  - $\sim x^{\text{exact}}(t) - X(t) = O(\Delta t^e)$
  - $\sim \Delta_{\text{conserved}} = O(\Delta t^e)$
- For Hamiltonian systems:  
symplectic vs non-symplectic

Higher orders ~~1~~: Verlet

Combine symplectic steps!

Semi-implicit Euler

V1

V2

$$\begin{cases} \vec{p}^{+1} = \vec{p} + \Delta t \vec{K} \\ \vec{x}^{+1} = \vec{x} + \Delta t \vec{p}^{+1} \end{cases}$$

$$\begin{cases} \vec{x}^{+1} = \vec{x} + \Delta t \vec{p} \\ \vec{p}^{+1} = \vec{p} + \Delta t \vec{K}^{+1} \end{cases}$$

Add 1/2 time step and alternate  $v^{1/2}$

Velocity-Verlet: (VV)

$$\vec{p}^{+1/2} = \vec{p} + \frac{1}{2} \Delta t \vec{K}$$

$$\vec{x}^{+1} = \vec{x} + \Delta t \vec{p}_{1/2}$$

$$\vec{p}^{+1} = \vec{p}_{1/2} + \frac{1}{2} \Delta t \vec{K}_{1/2}$$

Position Verlet: (PV)

$$\vec{x}^{+1/2} = \vec{x} + \frac{1}{2} \Delta t \vec{p}$$

$$\vec{p}^{+1} = \vec{p} + \Delta t \vec{K}_{1/2}$$

$$\vec{x}^{+1} = \vec{x}_{1/2} + \frac{1}{2} \Delta t \vec{p}^{+1}$$

Can also generalize it to higher order ( $\Delta t^{10}$  in CL!)

See Julia notebook

Staggered Leapfrog:

PV but don't stop at  $\vec{x}^{+1}$ !

$$\vec{x}^{+1/2} = \vec{x}^{-1/2} + \Delta t \vec{p}$$

$$\vec{p}^{+1} = \vec{p} + \Delta t \vec{K}_{1/2}$$

$\Delta$  Need to compare quantities at same times.

Higher orders  $\times 2$ : Runge-Kutta (RK)

Directly write higher order discret. of the ODE.

→ Not symplectic but works for non-Hamiltonian systems

(Verlet needs kicks to be momentum independent)

For instance: Modified Euler (RK2)

$$\vec{x}^{n+1} = \vec{x} + \frac{\Delta t}{2} (\vec{p}^{(1)} + \vec{p}^{(2)})$$

$$\vec{p}^{n+1} = \vec{p} + \frac{\Delta t}{2} (\vec{k}^{(1)} + \vec{k}^{(2)})$$

$$\vec{p}^{(1)} = \vec{p} \quad \vec{k}^{(1)} = \vec{K}(\vec{x}, \vec{p})$$

$$\vec{p}^{(2)} = \vec{p} + \Delta t \vec{k}^{(1)} \quad \vec{k}^{(2)} = \vec{K}(\vec{x} + \Delta t \vec{p}^{(1)}, \vec{p}^{(1)})$$

See julia notebook

RK4:

$$\vec{x}^{n+1} = \vec{x} + \frac{\Delta t}{6} (p^{(1)} + 2p^{(2)} + 2p^{(3)} + p^{(4)})$$

$$\vec{p}^{n+1} = \vec{p} + \frac{\Delta t}{6} (k_1^p + 2k_2^p + 2k_3^p + k_4^p)$$

$$p^{(1)} = \vec{p}$$

$$p^{(2)} = \vec{p} + \frac{\Delta t}{2} k_1^p$$

$$p^{(3)} = \vec{p} + \frac{\Delta t}{2} k_2^p$$

$$p^{(4)} = \vec{p} + \Delta t k_3^p$$

$$k_1^p = \vec{K}(\vec{x}, \vec{p})$$

$$k_2^p = \vec{K}(\vec{x} + \frac{\Delta t}{2} p^{(1)}, p^{(1)})$$

$$k_3^p = \vec{K}(\vec{x} + \frac{\Delta t}{2} p^{(2)}, p^{(2)})$$

$$k_4^p = \vec{K}(\vec{x} + \Delta t p^{(3)}, p^{(3)})$$

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- Summary:
- Time evolution through time-stepping
  - Physics of the problem informs on algorithm
  - Accuracy of solver matters

② (1+1) Wave equation:

Almost the same!

$$-\partial_t^2 \phi(x,t) + \partial_x^2 \phi(x,t) = 0$$

① Discretize in space (see lecture 1)

$$\partial_x^2 \phi(n,t) = \frac{\phi(n+1,t) + \phi(n-1,t) - 2\phi(n,t)}{\Delta x^2}$$

$$\equiv \Delta_x^2 \phi(n,t)$$

② Reduce to coupled 1<sup>st</sup> order ODE's

$$\partial_t \phi(n,t) \equiv \pi(n,t)$$

$$\Rightarrow \partial_t \pi(n,t) = \Delta_x^2 \phi(n,t)$$

$$\partial_t \phi(n,t) = \pi(n,t)$$

2N ODE's. Use the same methods  
↑  
as before!

number lattice points

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Comments: • Memory usage can

be an issue  $\rightarrow$  Verlet / leapfrog  
only 1 copy of fields

vs

RK, more than 1.

- CFL condition

$$\Delta t < \frac{1}{\sqrt{\text{dimension}}} \Delta x$$

Idea of proof: Take  $f(x,t) = e^{ikx - i\omega t}$

Gain ansatz:  $f(x, t + \frac{\Delta t}{2}) = G f(x,t)$

$$f(x, t + \Delta t) = G^2 f(x,t)$$

Position - Verlet:

$$f(x, t + \frac{\Delta t}{2}) = f(x,t) + \frac{\Delta t}{2} p(x,t)$$

$$p(x, t + \Delta t) = p(x,t) + \Delta t \Delta_x^r f(x,t)$$

$$f(x, t + \Delta t) = f(x, t + \frac{\Delta t}{2}) + \frac{\Delta t}{2} p(x,t)$$

Using:  $\Delta_x^r f(x,t) = \frac{(2 \cos(k\Delta x) - 1) f(x,t)}{\Delta x^2}$

find eq. for G.  $\frac{\Delta t^2}{\Delta x^2} < 1$  imposes

$|G| \leq 1 \Rightarrow$  stable.

